



# Multiferroic Ordering of Hexagonal Manganites

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## Multiferroic $RMnO_3$

### Multiferroics or Ferroelectromagnets:

Materials with simultaneous long-range (anti-) ferromagnetic, ferroelectric and/or ferroelastic order.

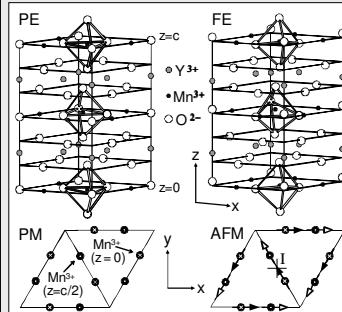
Hexagonal manganites  $RMnO_3$  ( $R = \text{Sc}, \text{Y}, \text{Ho}, \text{Er}, \text{Tm}, \text{Yb}, \text{Lu}$ )

$T < T_C \approx 600\text{-}1000 \text{ K} \Rightarrow$  ferroelectric (FEL) + paramagnetic (PM)

$T < T_N \approx 70\text{-}130 \text{ K} \Rightarrow$  ferroelectric (FEL) + antiferromagnetic (AFM)

$T < T_{RE} \approx 5 \text{ K} \Rightarrow$  FM or AFM order of  $R^{3+}$ -spins for  $R = \text{Ho} - \text{Yb}$

## Electric/magnetic order



Ferroelectric phase transition:  
Breaking of inversion symmetry II!

Order parameter: P

Antiferromagnetic phase transition of the Mn3+ sublattice:  
Breaking of time-reversal symmetry T, but not of inversion symmetry II!

Order parameter: ℓ

## Optical second harmonic generation (SHG)

In general: Multipole expansion of source term  $\vec{S}$  for SHG:

$$\vec{S} = \mu_0 \frac{\partial^2 \vec{P}^{NL}}{\partial t^2} + \mu_0 \left( \vec{\nabla} \times \frac{\partial \vec{M}^{NL}}{\partial t} \right) - \mu_0 \left( \vec{\nabla} \cdot \frac{\partial^2 \vec{Q}^{NL}}{\partial t^2} \right)$$

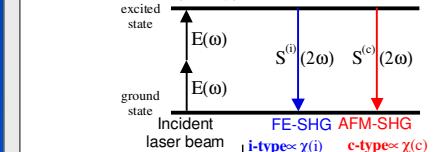
$\Rightarrow$  Three nonlinear contributions:

Electric dipole (ED):  $\vec{P}^{NL}(2\omega) \propto \hat{\chi}^{ED} : \vec{E}(\omega) \vec{E}(\omega)$

Magnetic dipole (MD):  $\vec{M}^{NL}(2\omega) \propto \hat{\chi}^{MD} : \vec{E}(\omega) \vec{E}(\omega)$

Electric quadrupole (EQ):  $\vec{Q}^{NL}(2\omega) \propto \hat{\chi}^{EQ} : \vec{E}(\omega) \vec{E}(\omega)$

SHG of electric-dipole type:



Source term for SHG:  $P_i(2\omega) = \epsilon_0 \chi_{ijk}^{SH} E_j(\omega) E_k(\omega)$

Intensity of SH signal:  $I_{SH} \propto |P(c) + P(i)|^2$

$$= (\chi^2(c) + A^2 \chi^2(i)) + 2A \chi(c) \chi(i) \cos \varphi I^2(\omega)$$

always  $> 0$       interference term

A: amplitude ratio of i-type and c-type terms  
 $\varphi$ : phase shift between complex contributions  
 A and  $\varphi$  can be fully controlled in experiment

Susceptibility  $\chi$  couples linearly to order-parameter!

## Observation of ferro-electromagnetic SHG

Order	Space group	Symmetry operation	Order parameter
PE + PM	$P6_3/mmc$	$I, T, IT$	---
FE + PM	$P6_3cm$	$T$	$P$
PE + AFM	$P6_3/mcm$	$I$	$\ell$
FE + AFM	$P6_3cm$	---	$P \cdot \ell$

Expansion of coupling of SHG to the order parameters:

$$\vec{P}^{NL}(2\omega) = \epsilon_0 (\hat{\chi}^{ED}(0) + \hat{\chi}^{ED}(\mathcal{P}) + \hat{\chi}^{ED}(\ell) + \hat{\chi}^{ED}(\mathcal{P}\ell)) \vec{E}(\omega) \vec{E}(\omega)$$

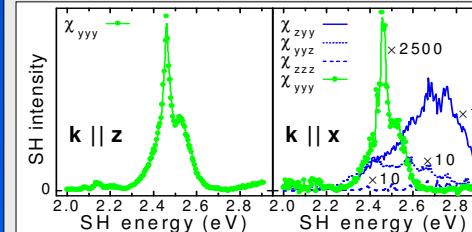
Lowest order non-zero contributions to SHG:

Zero order: Electric dipole (ED):  $\hat{\chi}^{ED}(P) = i_1, i_2, i_3$  and  $\hat{\chi}^{ED}(P \cdot \ell) = e_1$

First order: Magnetic dipole (MD):  $\hat{\chi}^{MD}(\ell) = m_1$

First order: Electric quadrupole (EQ):  $\hat{\chi}^{EQ}(\ell) = q_1, q_2, q_3$

	$S^{ED}(\mathcal{P})$	$S^{ED}(\mathcal{P} \cdot \ell)$	$S^{MD}(\ell)$	$S^{EQ}(\ell)$
$k \parallel x$	$S_y$	$2i_1 E_x E_z$	$e_1 E_y^2$	---
	$S_z$	$i_2 E_x^2 + i_3 E_z^2$	---	---
$k \parallel y$	$S_x$	$2i_1 E_x E_z$	---	$-2q_1 E_x$
	$S_z$	$i_2 E_x^2 + i_3 E_z^2$	$m_1 E_x$	$-q_2 E_x^2$
$k \parallel z$	$S_x$	---	$-2e_1 E_y E_z$	$-2q_3 E_x$
	$S_y$	---	$e_1 (E_x^2 - E_z^2)$	$q_3 (E_y^2 - E_z^2)$

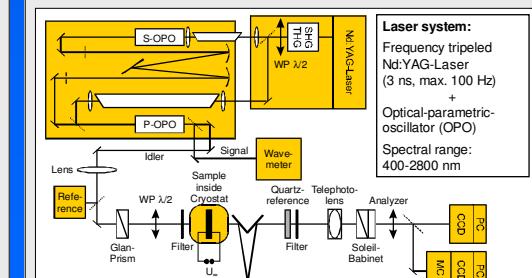


Identical magnetic spectra for  $k \parallel z$  and  $k \parallel x$

$\Rightarrow$  bilinear coupling to  $P \cdot \ell$ .

First observation of ferroelectromagnetic (FEM) SHG!

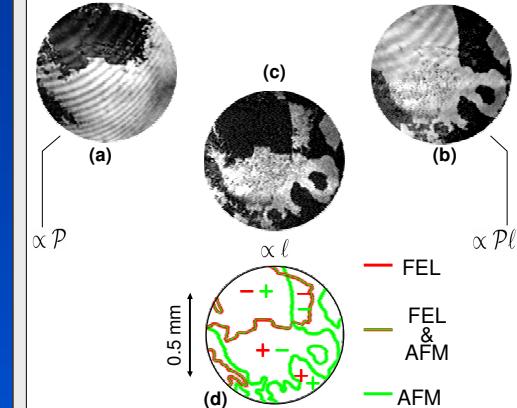
## Experimental setup



SHG/THG: Second/third harmonic generation, OPO: Optical-parametric oscillator, S/P-OPO: Seed/power OPO, WP: Waveplate, U: DC source-meter, MC: Monochromator, PM: Photomultiplier, CCD: Camera, PC: Computer

## Coexisting domains

Observed simultaneously in  $YMnO_3$  at 6 K



### Coexisting types of domains:

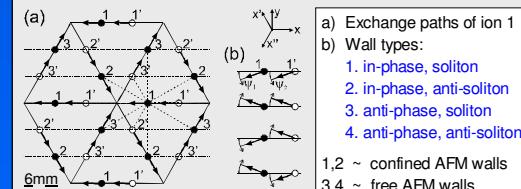
FEL:  $\propto P$    AFM:  $\propto \ell$    "Magneto-electric":  $\propto P\ell$

$P\ell = +1$  for  $P = \pm 1, \ell = \pm 1$     $P\ell = -1$  for  $P = \pm 1, \ell = \mp 1$

Any reversal of the FEL order parameter is clamped to a reversal of the AFM order parameter

"Free" and "confined" AFM walls

## Interaction of domain walls



a) Exchange paths of ion 1

b) Wall types:

1. in-phase, soliton

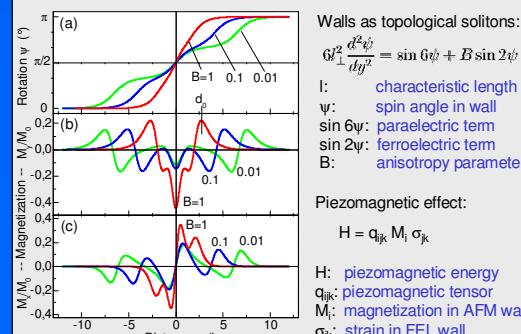
2. in-phase, anti-soliton

3. anti-phase, soliton

4. anti-phase, anti-soliton

1,2 = confined AFM walls

3,4 = free AFM walls



Walls as topological solitons:

$$6\frac{d^2\psi}{dy^2} = \sin 6\psi + B \sin 2\psi$$

I: characteristic length

ψ: spin angle in wall

sin 6ψ: paraelectric term

sin 2ψ: ferroelectric term

B: anisotropy parameter

Piezomagnetic effect:

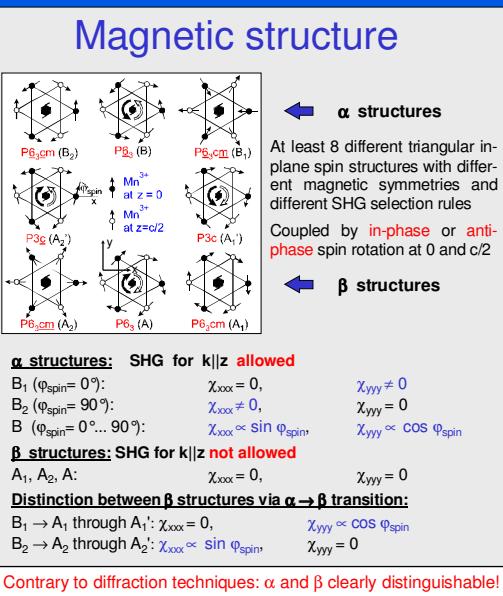
$$H = q_{jk} M_j \sigma_k$$

H: piezomagnetic energy

q\_{jk}: piezomagnetic tensor

M: magnetization in AFM wall

σ\_{jk}: strain in FEL wall



Contrary to diffraction techniques: α and β clearly distinguishable!