

Control of Magnetic Order by Electric Field in HoMnO₃

Th. Lottermoser⁴, Th. Lonkai¹, U. Amman^{1,2}, J. Ihringer¹, D. Hohlwein³, M. Fiebig⁴

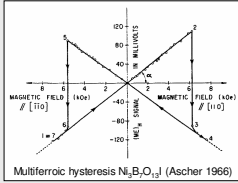
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¹Institut für Angewandte Physik, Universität Tübingen, ²Institut Laue-Langevin, Grenoble, ³Hahn-Meitner-Institut, Berlin, ⁴Max-Born-Institut, Berlin

Multiferroics

Compounds with simultaneous (anti-)ferromagnetic, ferroelectric and/or ferroelastic ordering

⇒ **Multiferroics** (Aizu 1969)



1958 Idea of new compounds with coexisting magnetic and electric ordering by Smolenskii and Ioffe

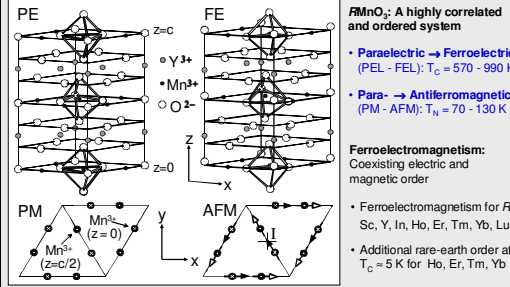
1966 First experimental proof of a "multiferroic effect" by Ascher et al.

1975 Suggestions for technical applications based on multiferroic properties by Wood and Austin

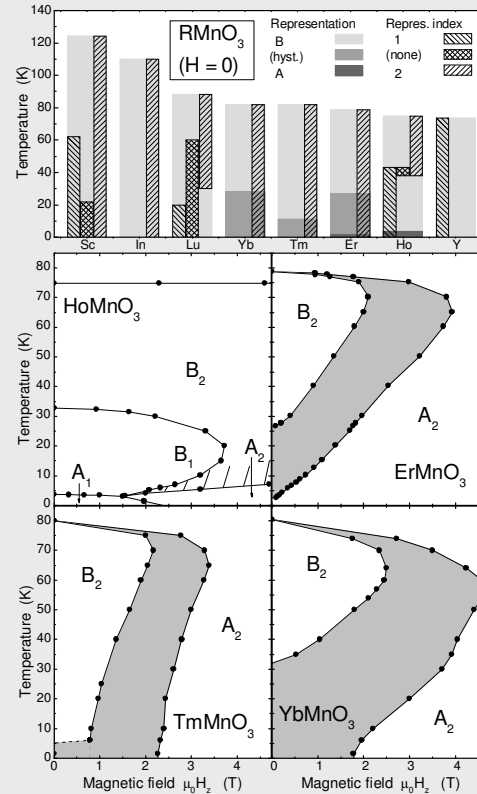
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2000 "Why are there so few magnetic ferroelectrics?" by Hill

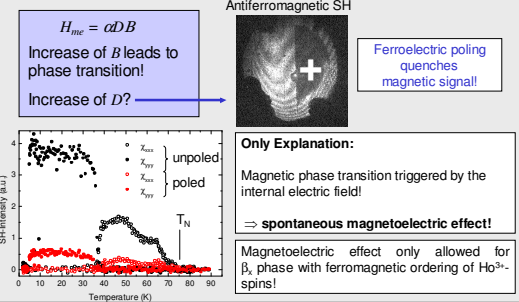
Electric and magnetic ordering of RMnO₃



H/T phase diagram



Electric field induced magnetoelectric effect



Linear magnetoelectric effect

Polarization and magnetization of a medium:

$$P_i = \epsilon_0 \chi_{ij}^e E_j, \quad M_i = \chi_{ij}^m H_j$$

Covariant relativistic formulation:

$$\mu_{\alpha\beta} c M^{\alpha\beta} = \frac{1}{2} \epsilon_0 \epsilon^{\alpha\beta} F^{\mu\nu} \text{with:}$$

$$F_{\mu\nu} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -cB_z & cB_y \\ E_y & cB_z & 0 & -cB_x \\ E_z & -cB_y & cB_x & 0 \end{pmatrix}$$

Relativistic equivalence of electric and magnetic fields requires "magneto-electric" cross-correlation ($\sim \alpha$) in matter:

$$P_i = \epsilon_0 \chi_{ij}^e E_j + \frac{1}{c} \alpha_{ij} H_j, \quad M_i = \chi_{ij}^m H_j + \frac{1}{\mu_0 c} \alpha_{ji} E_j$$

1960: Small effect (10⁻⁵)
Limited choice of compounds
Theoretically not well understood

2000: New materials with structural (heterostructures) or gigantic (multiferroic) magnetoelectric effects
New theoretical concepts

Second harmonic generation

SHG of electric-dipole type:

excited state
ground state
Incident laser beam
FE-SHG AFM-SHG
i-type $\propto \chi(i)$ c-type $\propto \chi(c)$

Superposition

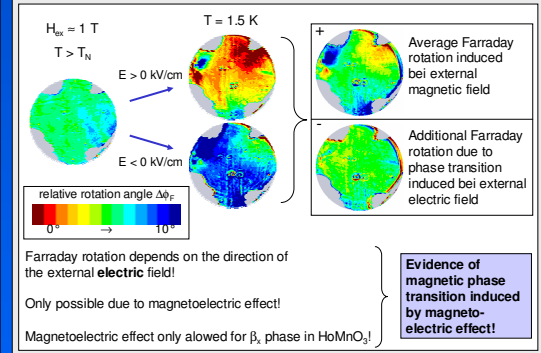
Source term for SHG: $P_i(2\omega) = \epsilon_0 \chi_{ijk}^{SH} E_j(\omega) E_k(\omega)$

Intensity of SH signal: $I_{SH} \propto |P(c) + P(i)|^2 \propto |\chi(c) + A e^{i\varphi} \chi(i)|^2 I^2(\omega) = (\chi^2(c) + A^2 \chi^2(i) + 2A \chi(c) \chi(i) \cos \varphi) I^2(\omega)$

always > 0 interference term

A: amplitude ratio of i-type and c-type terms
 φ : phase shift between complex contributions
A and φ can be fully controlled in experiment

Magnetization control by electric field



Multiferroic RMnO₃

Manganites RMnO₃ (R = Sc, Y, In, Ho, Er, Tm, Yb, Lu):

Multiferroic compounds with simultaneous ferroelectric and antiferromagnetic ordering.

Question: Results the multiferroic nature in the existence of any magnetoelectric effects?

Experimental methods:

- Second harmonic generation (SHG)
 - Polarization dependent spectroscopy
 - Nonlinear phase sensitive domain topography
- Neutron diffraction

Magnetic structure and SHG selection rules

At least 8 different triangular in-plane spin structures with different magnetic symmetries and different selection rules for SHG

α structures: SHG for k||z allowed

α_x ($\varphi = 0^\circ$): $\chi_{xxx} = 0, \chi_{yyy} \neq 0$
 α_y ($\varphi = 90^\circ$): $\chi_{xxx} \neq 0, \chi_{yyy} = 0$
 α_β ($\varphi = 0-90^\circ$): $\chi_{xxx} \propto \sin \varphi, \chi_{yyy} \propto \cos \varphi$

β structures: SHG for k||z not allowed

$\beta_x, \beta_y, \beta_z$: $\chi_{xxx} = 0, \chi_{yyy} = 0$

Determine β structure from α - β transition

$\alpha_x \rightarrow \beta_y$: $\chi_{xxx} = 0, \chi_{yyy} \propto \cos \varphi$
 $\alpha_x \rightarrow \beta_z$: $\chi_{xxx} \propto \sin \varphi, \chi_{yyy} = 0$

Contrary to diffraction techniques: α and β models clearly distinguishable!

Magnetoelectric 3d-4f superexchange in RMnO₃

Gigantic magnetoelectric effect which originates in 3d-4f superexchange; triggers hidden phase transition!

$H_{ex} = \sum_{k=1}^{4(4-3)} \sum_{l=1}^6 \sum_{j=1}^6 \tilde{S}^{kl}(a) \tilde{A}^{kl,j} \tilde{S}^{jkl}(a)$

k: R sites with 3 and 3m symmetries
l: all R ions at k sites (4+2)
j: 6 Mn ions neighboring an R ion
A: Mn-R exchange matrix (4 types)
S: spins of Mn and R ions

α model: $H_{ex}(\alpha) = 0$ no change!
 β model: lowers ground-state energy:
 $H_{ex}(\beta_x) = 6t S^R S^{Mn} [(A_{12}^{Mn} - A_{12}^{Mn}) - (A_{13}^M - A_{13}^M)]$
 $H_{ex}(\beta_y) = 6t S^R S^{Mn} [(A_{12}^{Mn} + A_{12}^{Mn}) - (A_{13}^M + A_{13}^M)]$

Ferroelectric distortion modifies the superexchange:
 $\delta A \equiv A - \Delta$, $\delta A = \delta A_{ij} P_i P_j$ Scales with order par.
 $\alpha_{12} \equiv 6t S^{Mn} (\delta A_{12}^{Mn} \pm \delta A_{12}^{Mn})_{12}$ Substitution leads to:
 $H_{ex}(\beta_x) = \alpha_{12} P_x S^R$ **ME contribution**

Neutron diffraction results

