

# Determination of the Magnetic Symmetry of Hexagonal Manganites by Second Harmonic Generation

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## Motivation

In the hexagonal manganites  $\text{RMnO}_3$  ( $R = \text{Y, Sc, Ho, Er, Tm, Yb, Lu}$ ), a phase transition from the paramagnetic into the antiferromagnetic phase occurs at  $T_N = 70 - 130$  K. With neutron or magnetic x-ray diffraction experiments it is *not* possible to determine unambiguously the antiferromagnetic structure of the frustrated  $\text{Mn}^{3+}$  spins.

⇒ Introduction of a new optical method:

Polarisation dependent spectroscopy of magnetic second harmonic (MSHG)

## Method

Second-order nonlinear polarisation:

$$P_i(2\omega) = \epsilon_0 \sum_{j,k=1}^3 \chi_{ijk}^{SH} E_j(\omega) E_k(\omega)$$

Second-order nonlinear susceptibility in magnetically ordered crystals:

$$\chi_{ijk}^{SH} = \chi_{ijk}^{cr}(i) + \chi_{ijk}^{mag}(c)$$

determined by crystal structure

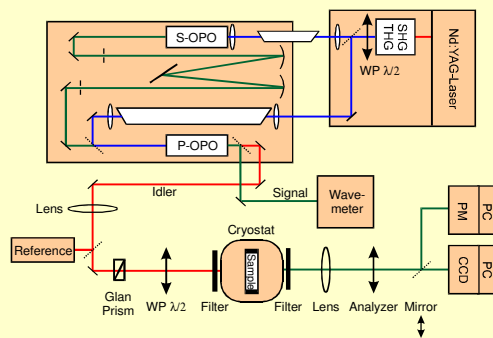
determined by magnetic structure

Studying the structure of the nonlinear susceptibility  $\chi_{ijk}^{mag}(c)$  allows the determination of the magnetic structure. The nonvanishing components of  $\chi_{ijk}^{mag}(c)$  are given by measurements of the

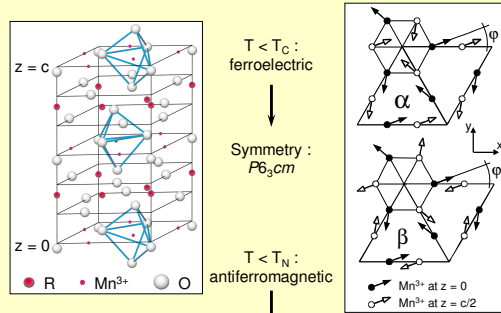
spectral, polarization, and temperature dependence.

Spatially resolved measurements lead to additional information.

## Experimental setup



## Crystal and spin structure



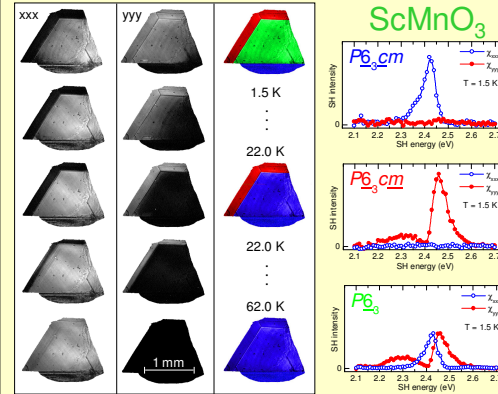
Some magnetic symmetry and SH selection rules:

$\alpha_1$ ( $\varphi = 0^\circ$ ): $P6_3cm$	$\beta_1$ ( $\varphi = 0^\circ$ ): $P6_3cm$
$\chi_{yyy} = -\chi_{yxx} = -\chi_{xyx} = -\chi_{xxy}$	$\chi_{xyz} = \chi_{xzy} = -\chi_{yxz} = -\chi_{zyx}$
$\alpha_2$ ( $\varphi = 90^\circ$ ): $P6_3cm$	$\beta_2$ ( $\varphi = 90^\circ$ ): $P6_3cm$
$\chi_{xxx} = -\chi_{xyy} = -\chi_{xyx} = -\chi_{yxx}$	$\chi_{zxx} = \chi_{zyy}, \chi_{xxz} = \chi_{yyz} = \chi_{yzy}, \chi_{zzz}$
$\alpha_\varphi$ ( $\varphi = 0^\circ \dots 90^\circ$ ): $P6_3$	$\beta_\varphi$ ( $\varphi = 0^\circ \dots 90^\circ$ ): $P6_3$
$\chi(\alpha_1) \otimes \chi(\alpha_2)$	$\chi(\beta_1) \otimes \chi(\beta_2)$

⇒ MSHG allowed for  $k \parallel z$ !

⇒ No MSHG allowed for  $k \parallel z$ !

## Spin-angle topography



Magnetic symmetry:

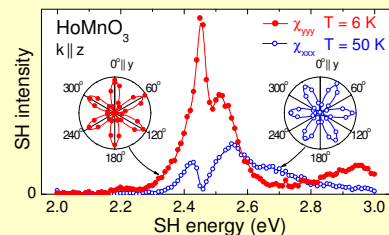
$T < T_N$ :  
 $P6_3cm$   
 $T < 62$  K:  
 $P6_3cm + P6_3cm$   
 $T < 22$  K:  
 $P6_3cm + P6_3cm + P6_3$

Spin-angle

$E_{SH} = 2.44$  eV  
 $T = 1.5$  K

$$\varphi = \arctan\left(\frac{\chi_{yyy}^{mag}(c)}{\chi_{xxx}^{mag}(c)} \sqrt{I_{xxx} / I_{yyy}}\right)$$

## Phase transitions in HoMnO3 and LuMnO3



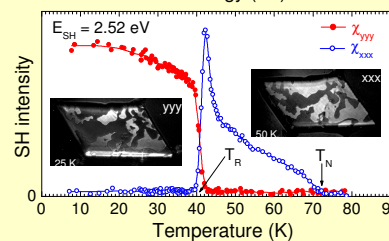
HoMnO<sub>3</sub>:

Additional magnetic phase transition at 41 K:

$T < T_R$ : Symmetry  $P6_3cm$

$T > T_R$ : Symmetry  $P6_3cm$

90° spin rotation at  $T_R = 41$  K

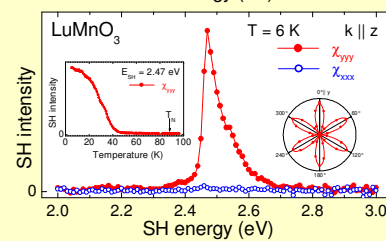
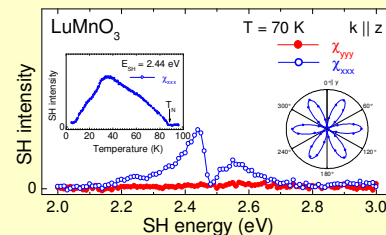


LuMnO<sub>3</sub>:

Temperature dependence of the symmetry:

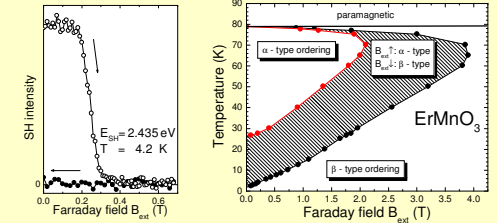
$T < T_N$ : Symmetry  $P6_3cm$

$T \rightarrow 0$  K: Transfer from  $P6_3cm$  to  $P6_3cm$  symmetry

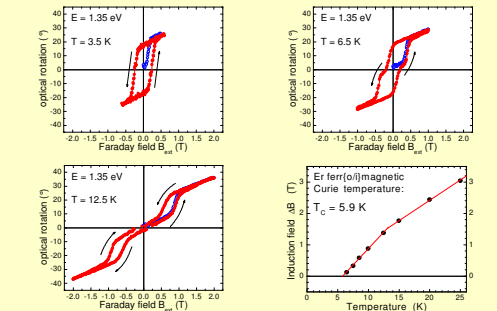


## Interaction of Mn and Er magnetic sublattices

Quenching of Mn-induced SH signal due to an antiferromagnetic transition from  $\alpha$ - to  $\beta$ -type ordering



Linear optical rotation induced by Er ferro- or ferrimagnetic moment



## Conclusions

- MSHG is an important complement to neutron and x-ray diffraction experiments for symmetry analysis
- Additional degrees of freedom:
  - Spectroscopy
  - Topography
- $\alpha$ -type ordering of hexagonal  $\text{RMnO}_3$  compounds
- Phase coexistence  $\Rightarrow$  small in plane anisotropy

Symmetry of the hexagonal manganites  $\text{RMnO}_3$

